

A half de Sitter Holography

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[arXiv:2306.07575](https://arxiv.org/abs/2306.07575) with Taishi Kawamoto, Yu-ki Suzuki, Tadashi Takayanagi

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01. Introduction- de Sitter Space

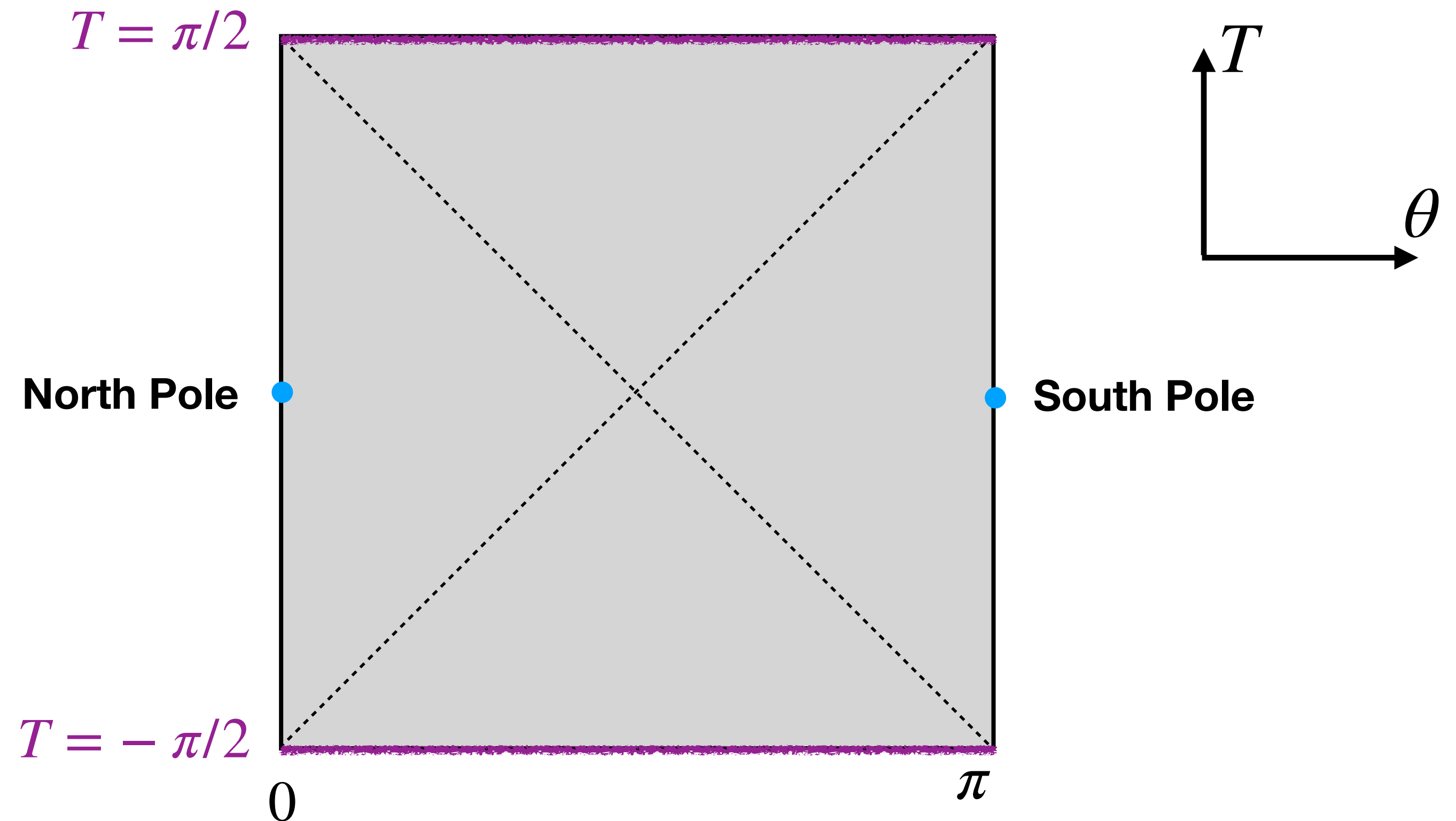
dS/CFT Correspondence

[Strominger 2001]

$$ds^2 = L^2 \left(-dt^2 + \cosh^2 t (d\theta^2 + \sin^2 \theta d\Omega_{d-1}^2) \right)$$

$$= \frac{L^2}{\cos^2 T} \left(-dT^2 + d\theta^2 + \sin^2 \theta d\Omega_{d-1}^2 \right)$$

$$S_{\text{dS}} = \frac{\text{Area}}{4G_{\text{N}}} = \frac{L^{d-1} \Omega_{d-1}}{4G_{\text{N}}} \equiv N$$

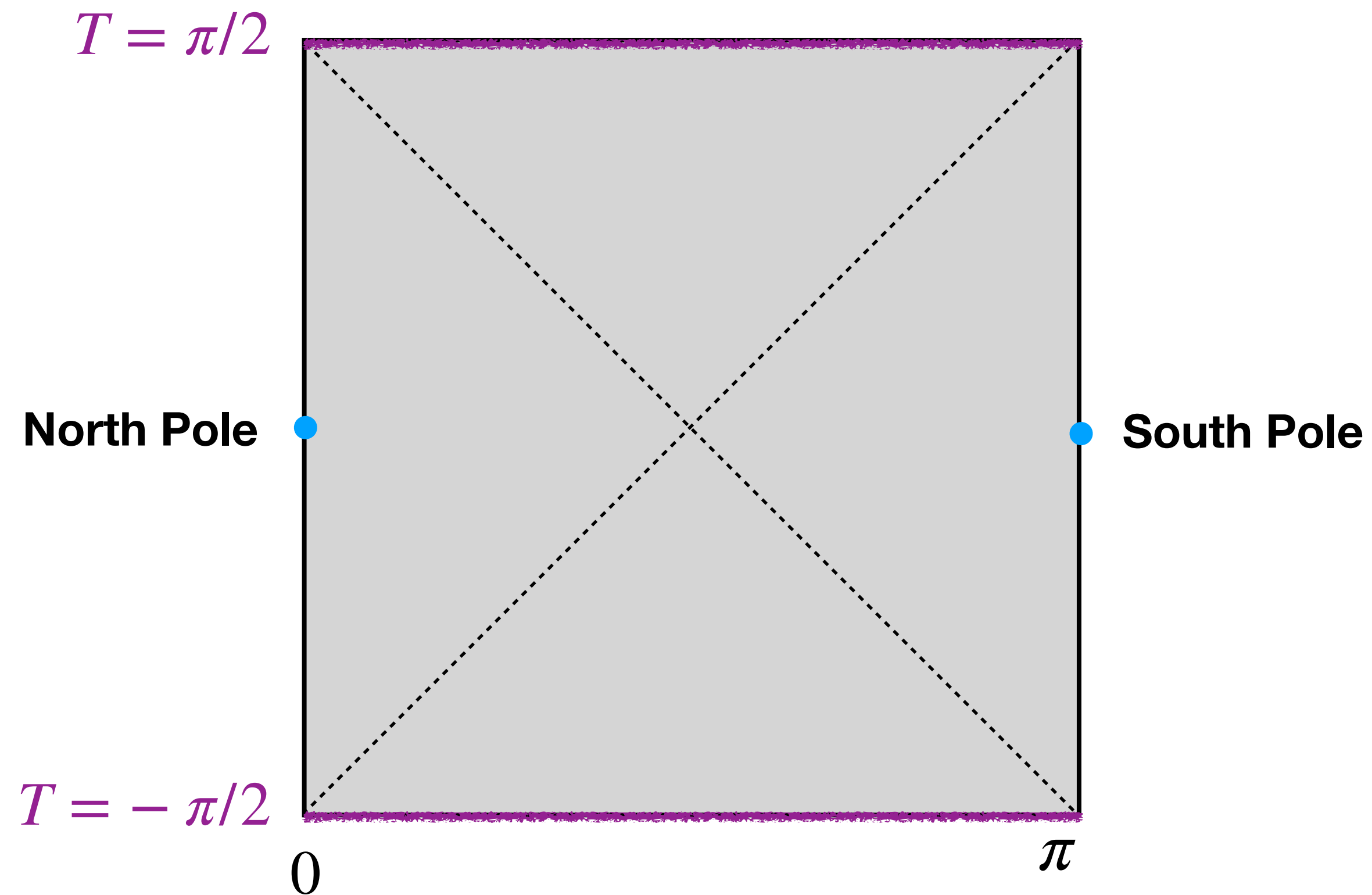


“Analytical continuation” of AdS/CFT correspondence

dS/CFT correspondence: gravity in asymptotically dS_{d+1} spacetime
 a d -dimensional CFT living on the spacelike boundary at **(timelike)** future infinity of dS.

dS/CFT Correspondence

[Strominger 2001]



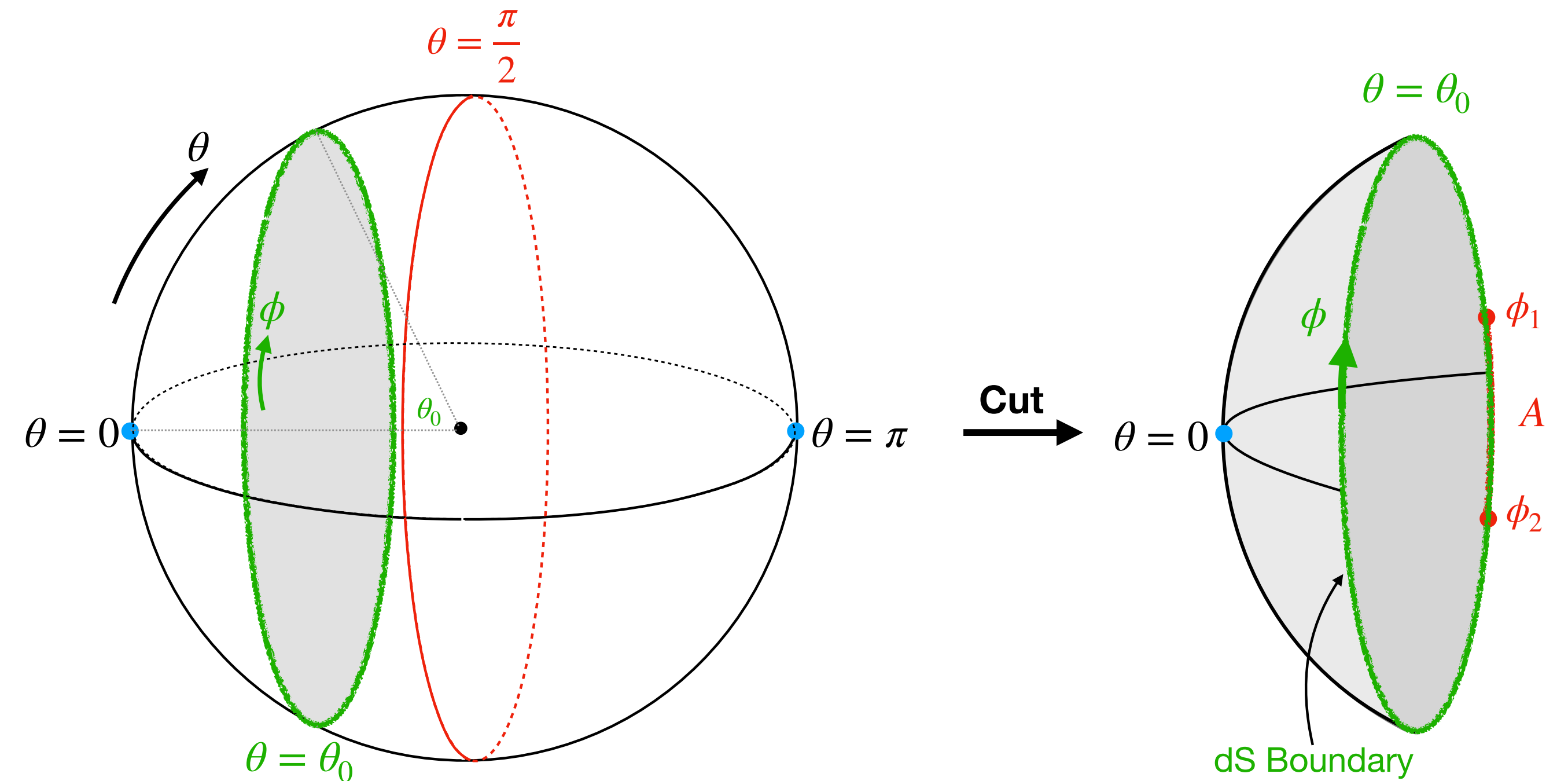
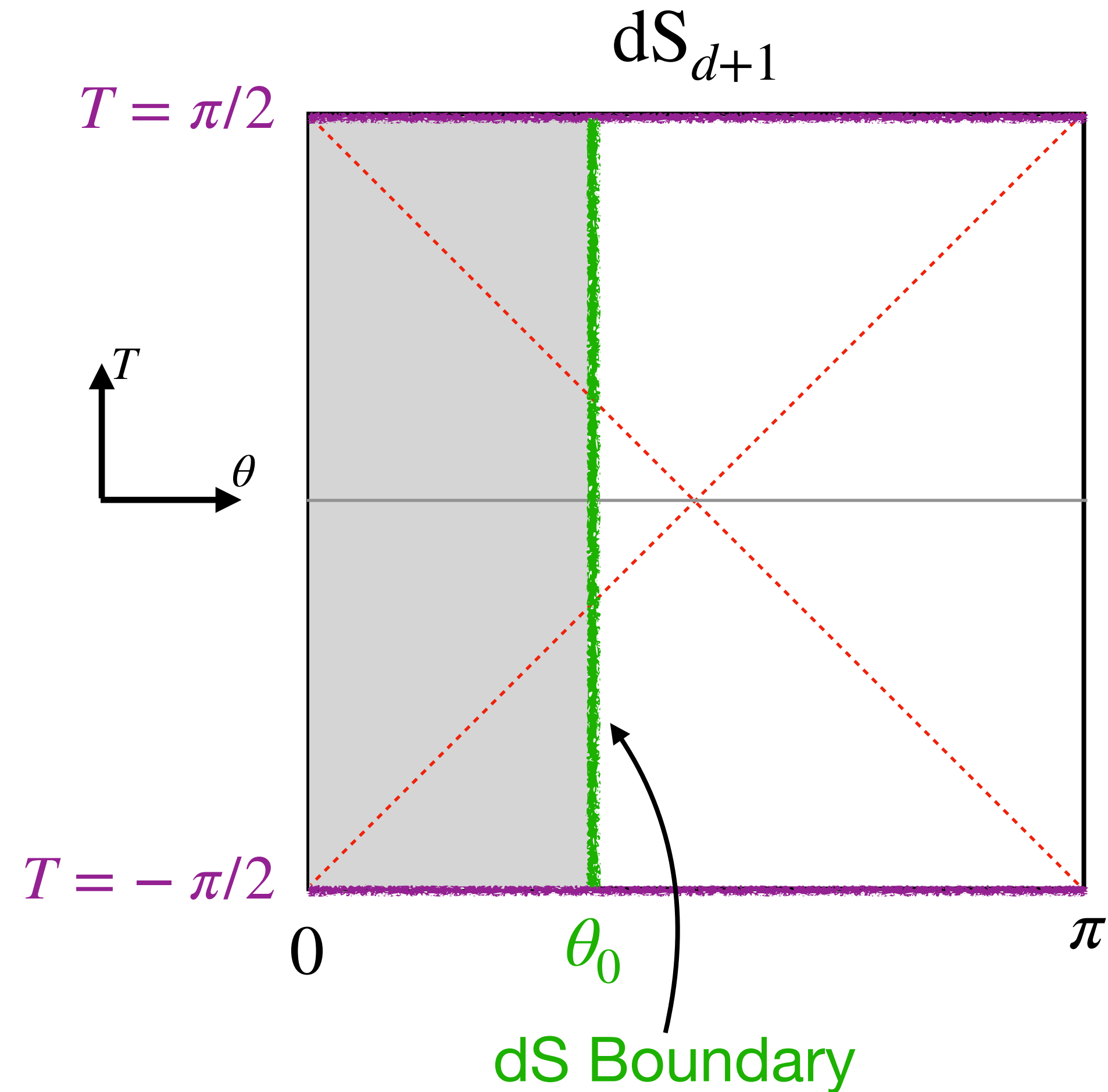
dS₄ [Anninos-Hartman-Strominger 2011]

dS₃ [Hikida,Nishioka,Takayanagi,Taki 2021]

- de Sitter is forbidden as supersymmetry solution to supergravity
- Dual CFT in dS/CFT is exotic, e.g., dS₃/CFT₂, pure imaginary central charge

$$c = i \frac{3L_{\text{dS}}}{2G_N}$$

Topic Today: A half de Sitter space



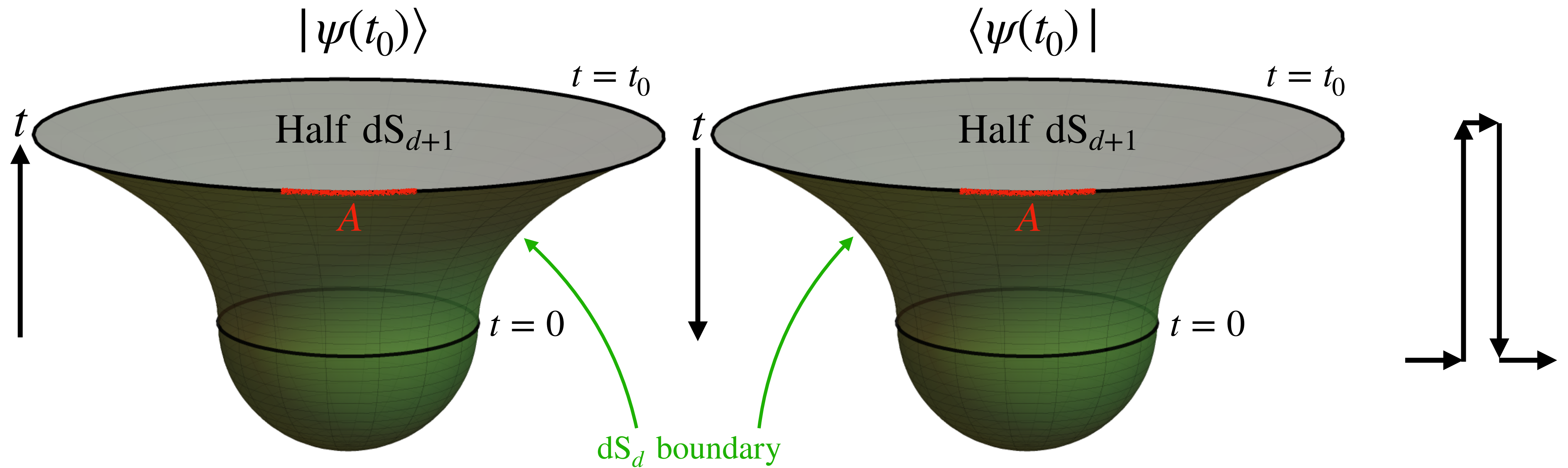
$$ds^2 = -dt^2 + \cosh^2 t (d\theta^2 + \sin^2 \theta d\Omega_{d-1}^2)$$

d+1-dim de Sitter Bulk

$$ds^2 = -dt^2 + \cosh^2 t \sin^2 \theta_0 d\Omega_{d-1}^2$$

d-dim de Sitter Boundary

Perspective 01: Schwinger-Keldysh Prescription



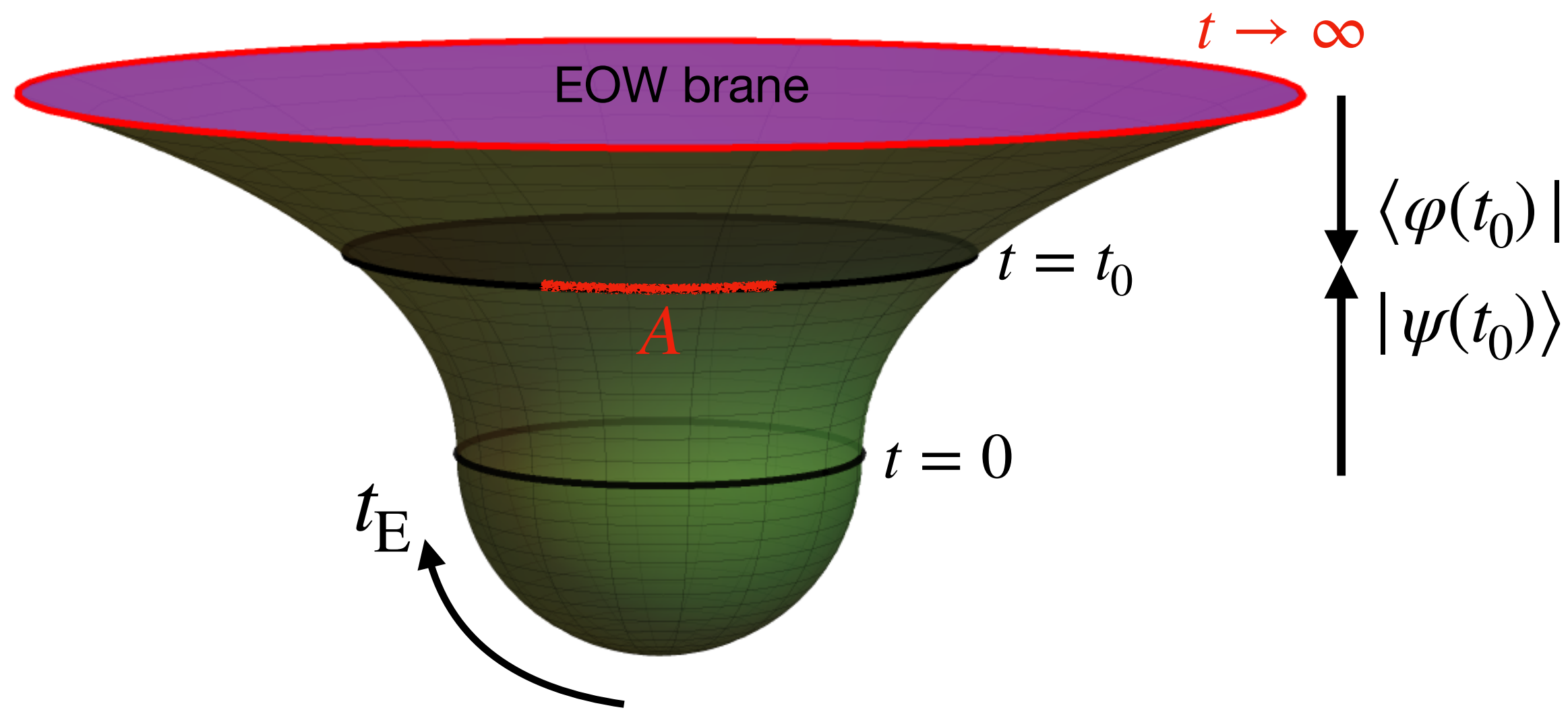
Holographic Entanglement Entropy

$$\rho_A = \text{Tr}_B \left[\frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle} \right]$$

$$S_A = \text{Tr} [-\rho_A \log \rho_A]$$

Perspective 02: Final State Projection

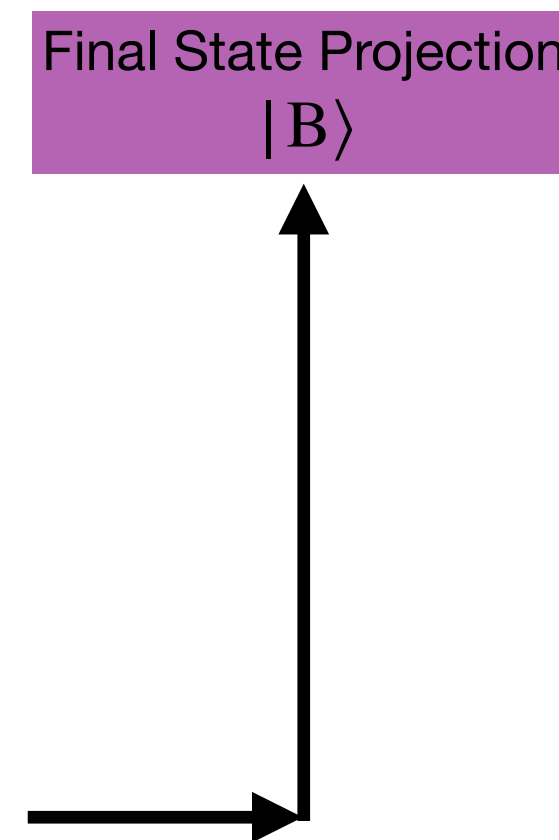
Final state projection at future infinity



Reduced transition matrix

$$\tau_A = \text{Tr}_B \left[\frac{|\psi\rangle \langle \varphi|}{\langle \psi | \varphi \rangle} \right]$$

It is not always Hermitian!



A generalization of entanglement entropy

Pseudo entropy $S_A = \text{Tr} [-\tau_A \log \tau_A]$

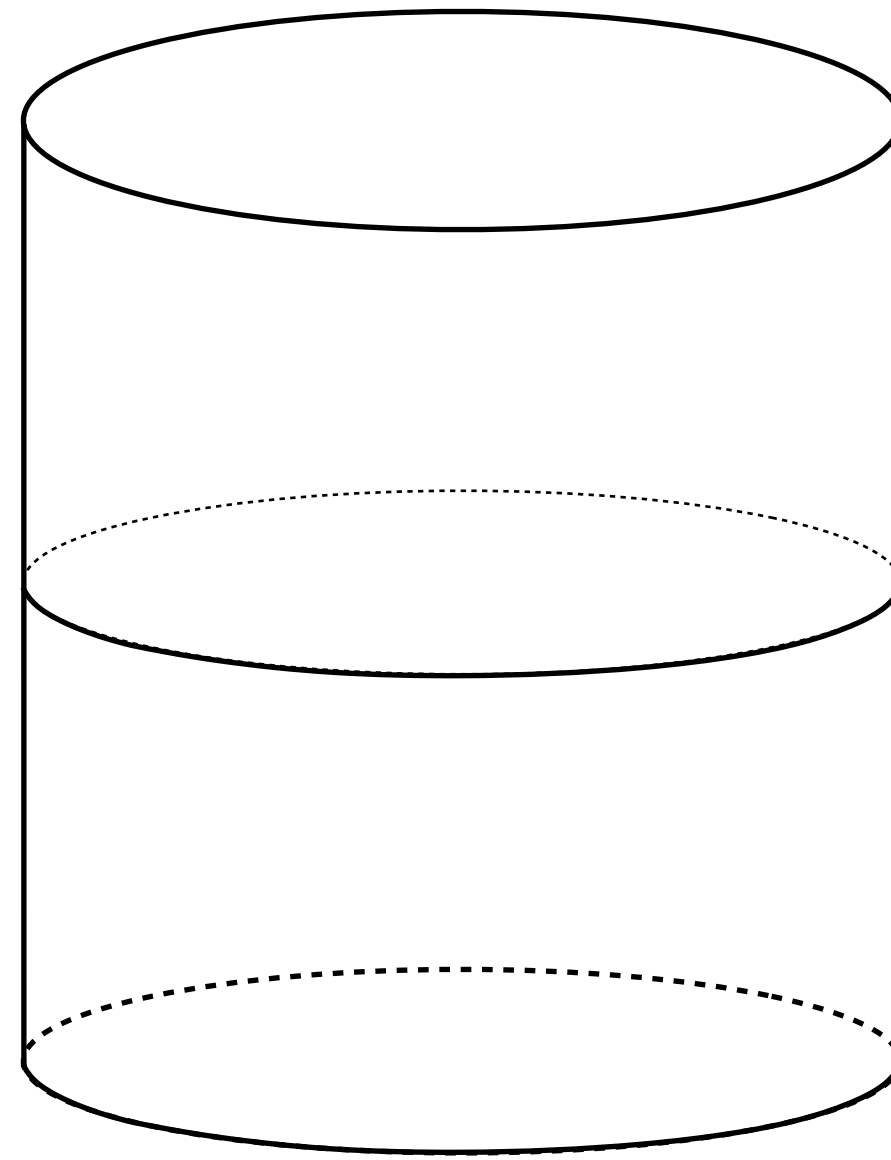
Complex Values!

Holographic Pseudo Entropy

[Nakata, Takayanagi, Taki, Tamaoka, Wei 2020]

Note: Its holographic dual is also given by the area of the minimal surface

02. CFT on dS space via AdS/CFT



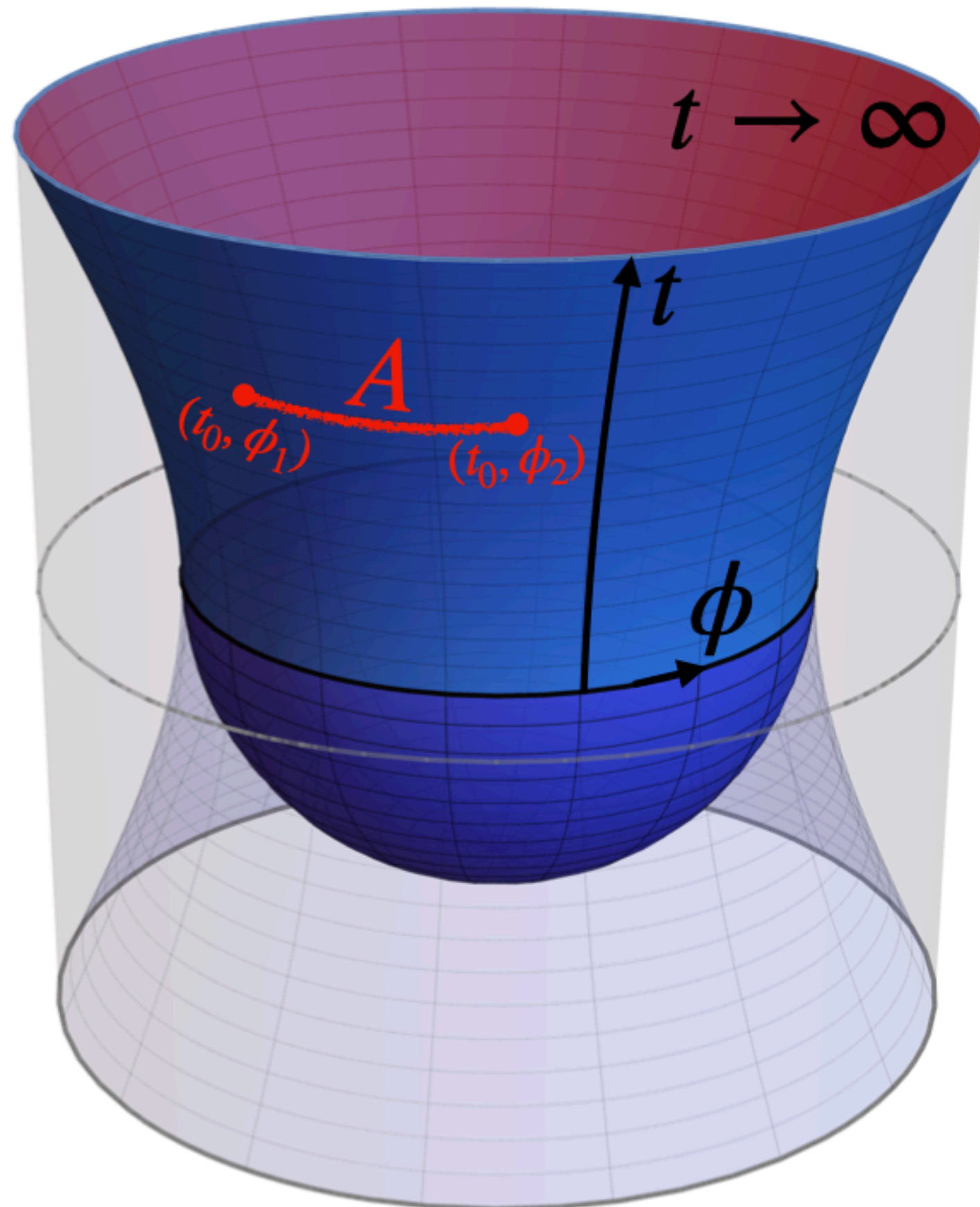
Perspective 01: Holographic Entanglement Entropy

Hartle-Hawking State
“No boundary proposal”

[“Wave Function of the Universe” 1983]

Lorentzian
Evolution

Euclidean
Evolution



AdS₃/CFT₂

$$S_A^{\text{con}} = \frac{c}{3} \log \left[\frac{2 \cosh t_0 \sin \frac{|\phi_1 - \phi_2|}{2}}{\epsilon} \right],$$

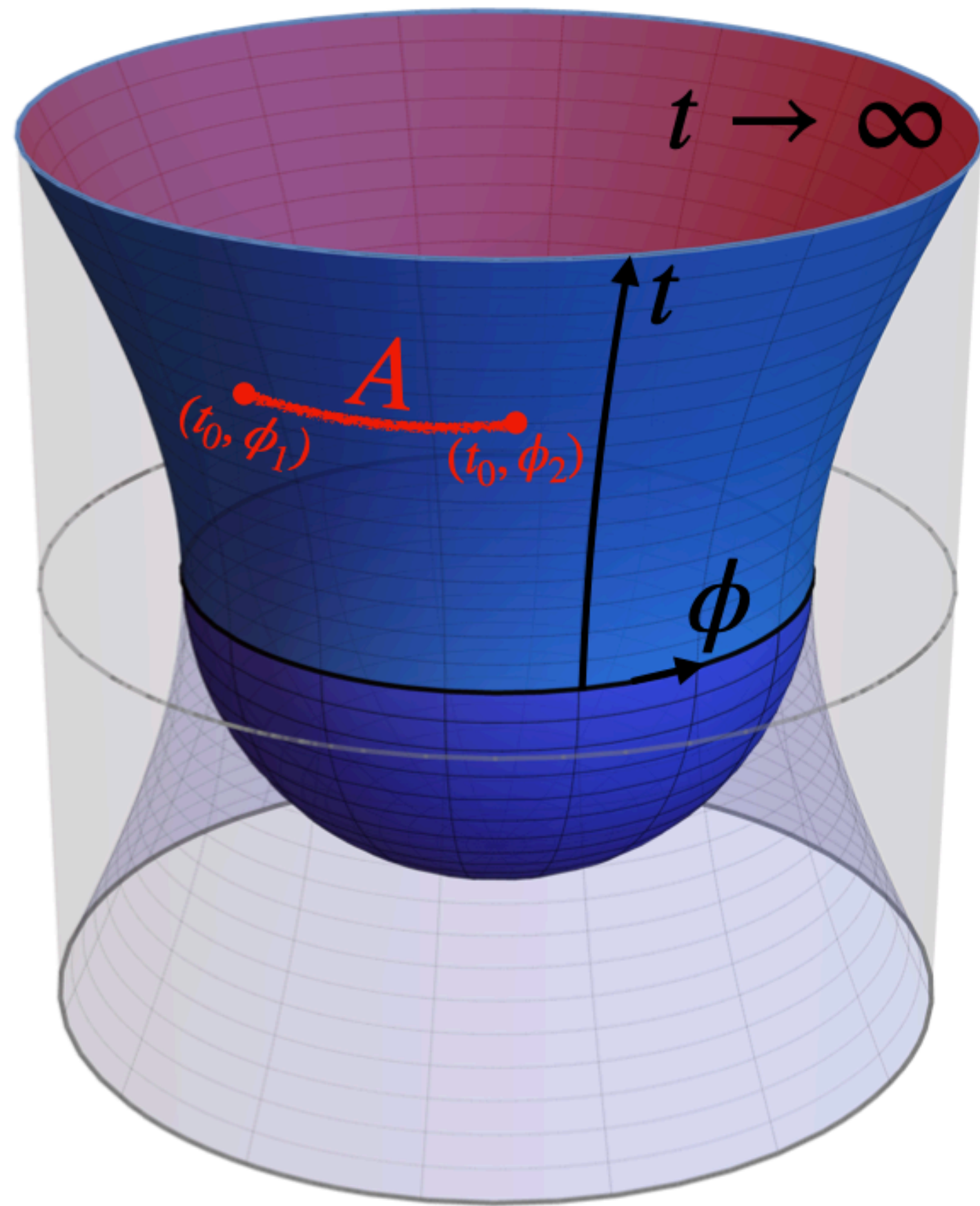
Linear growth at late times

Schwinger-Keldysh prescription

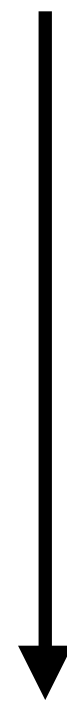
Future infinity plays no role!

Perspective 02: Holographic Pseudo Entropy

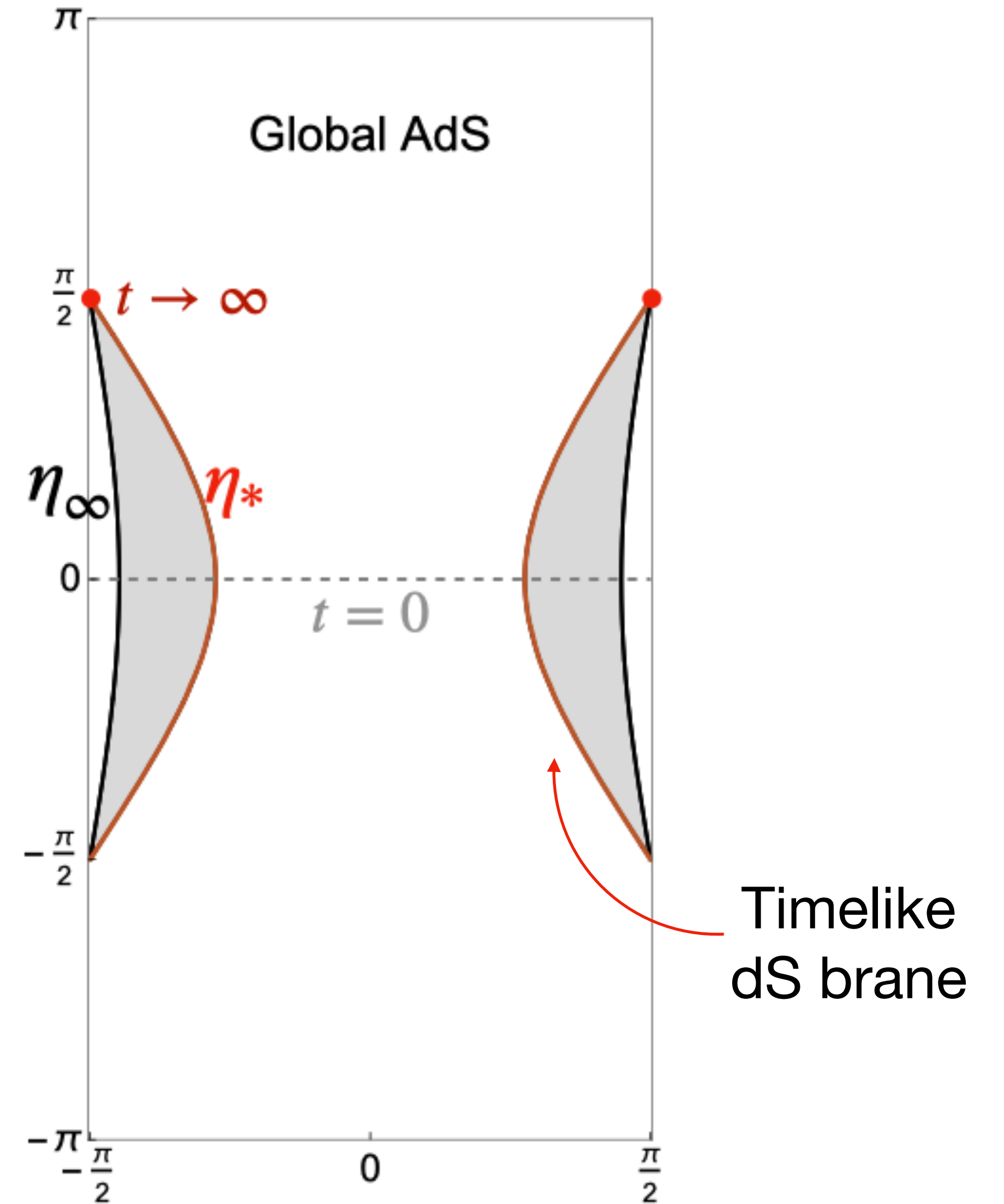
Boundary state



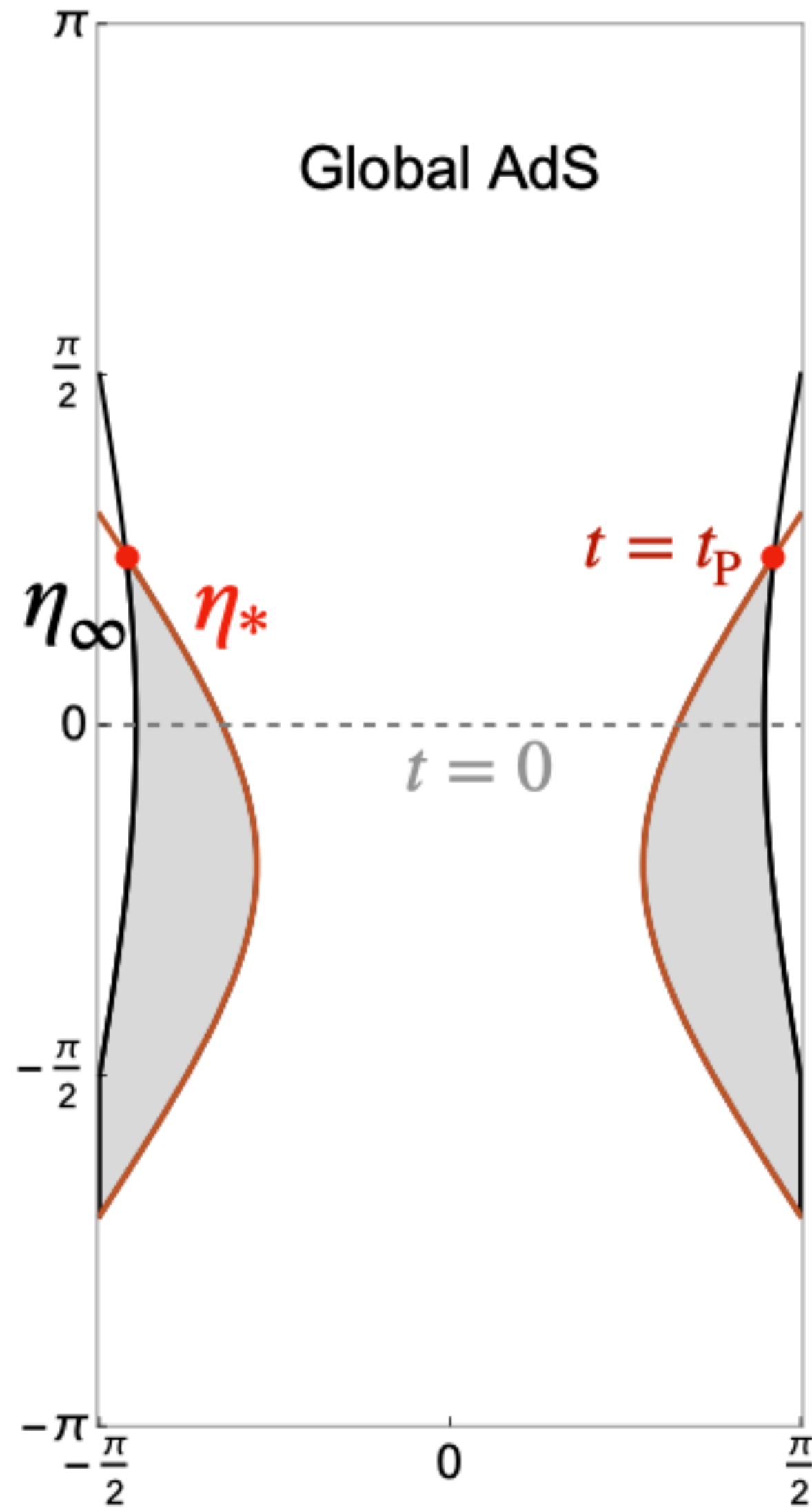
Final state projection
on the boundary



End-of-the-World brane
in the bulk



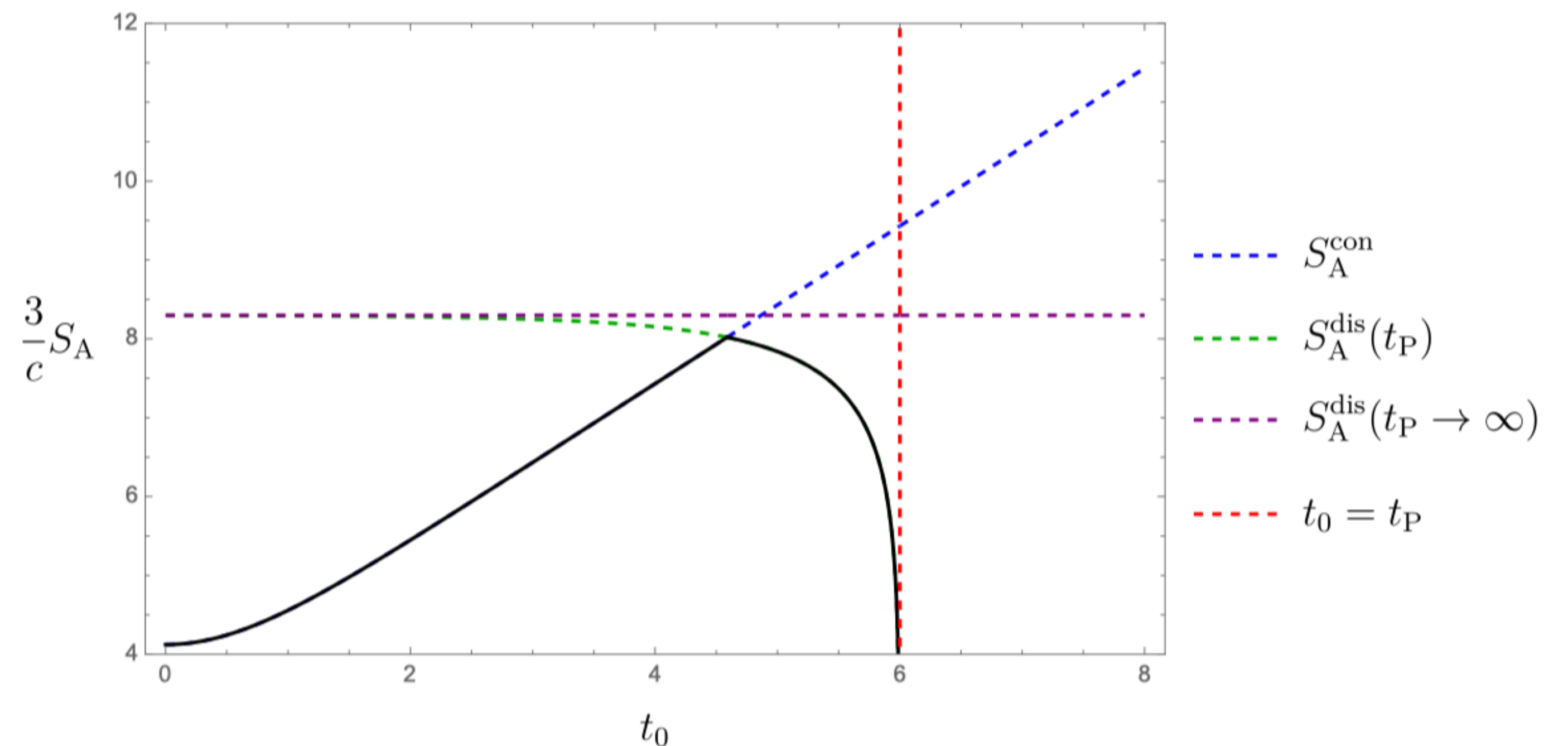
Perspective 02: Holographic Pseudo Entropy



Final State Projection at a finite time

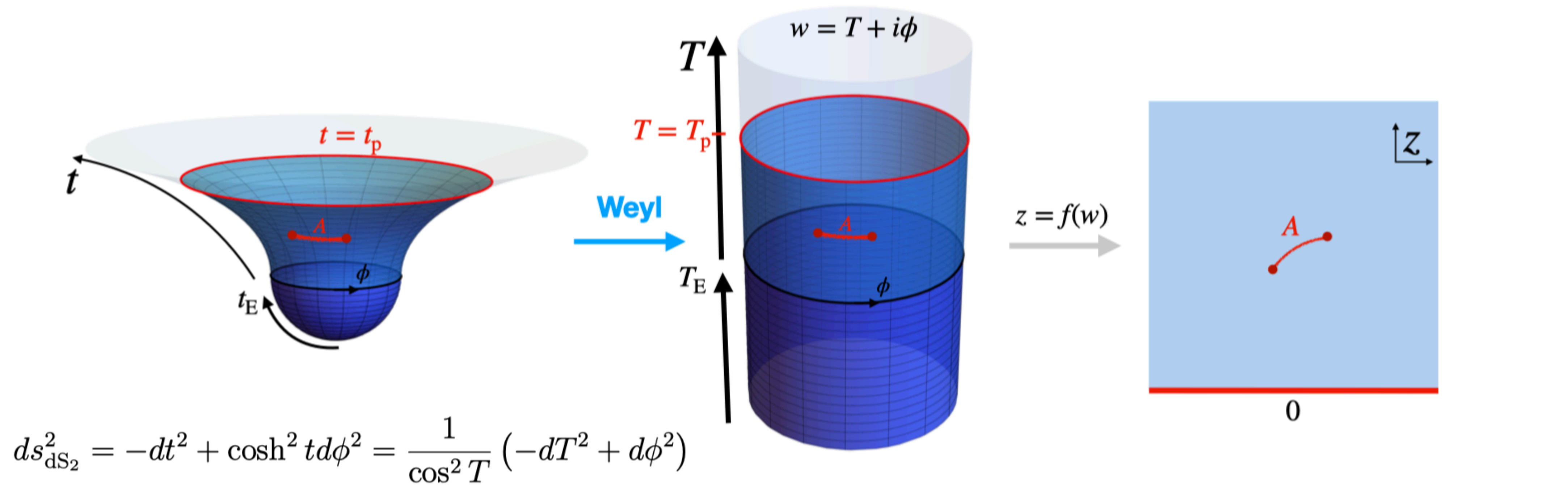
$$S_A^{\text{dis}} = \frac{c}{3} \log \left[\frac{2}{\epsilon} \cdot \frac{\sinh t_P - \sinh t_0}{\cosh t_P} \right] - \frac{c}{3} \eta_*$$

Boundary entropy



Perspective 02: Holographic Pseudo Entropy

Calculations from CFT_2 on dS_2

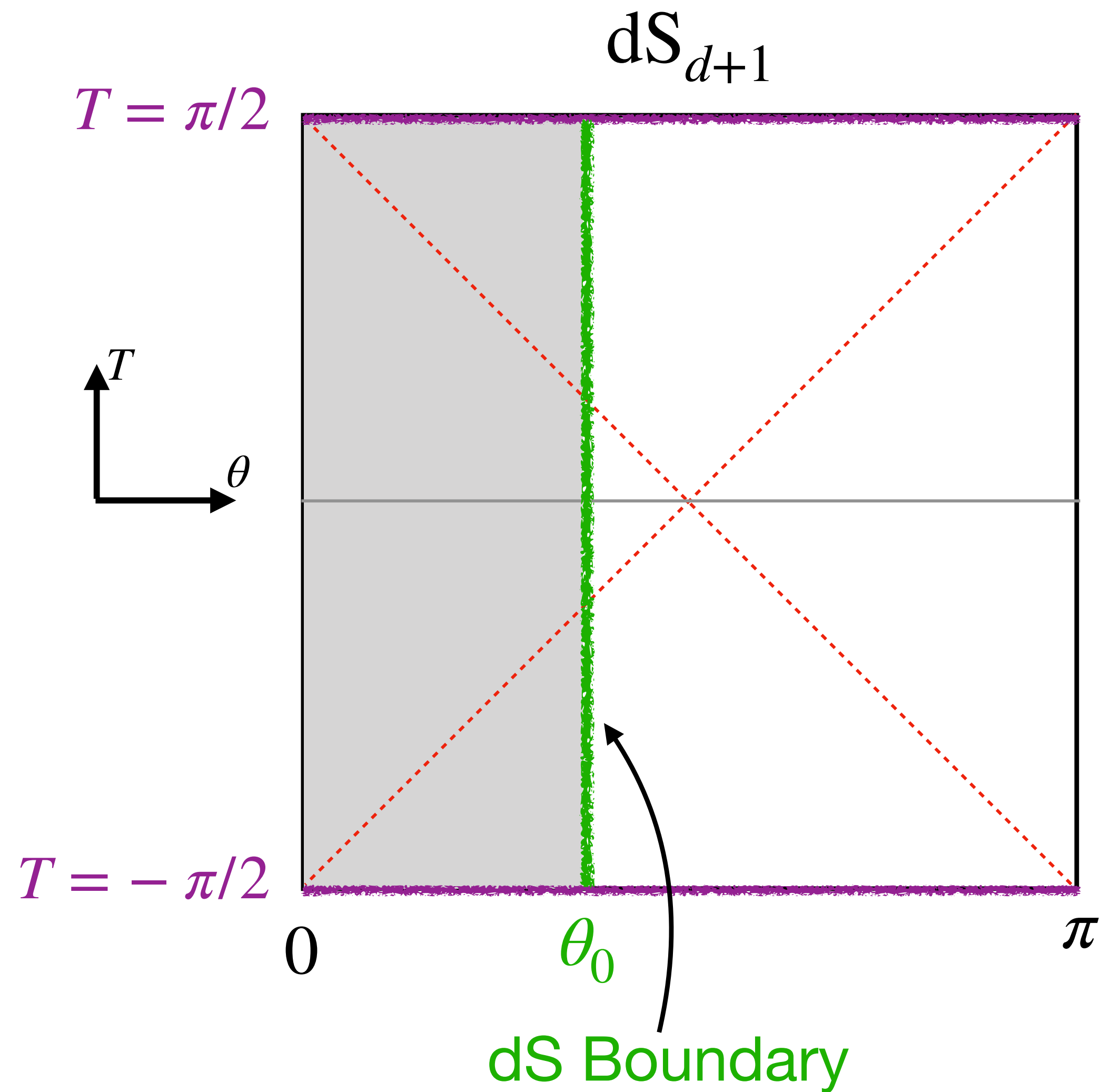


$$S_A^{\text{con,cyl}} = \frac{c}{6} \log \frac{|f(w_1) - f(w_2)|^2}{\epsilon^2 |f'(w_1)| |f'(w_2)|}$$

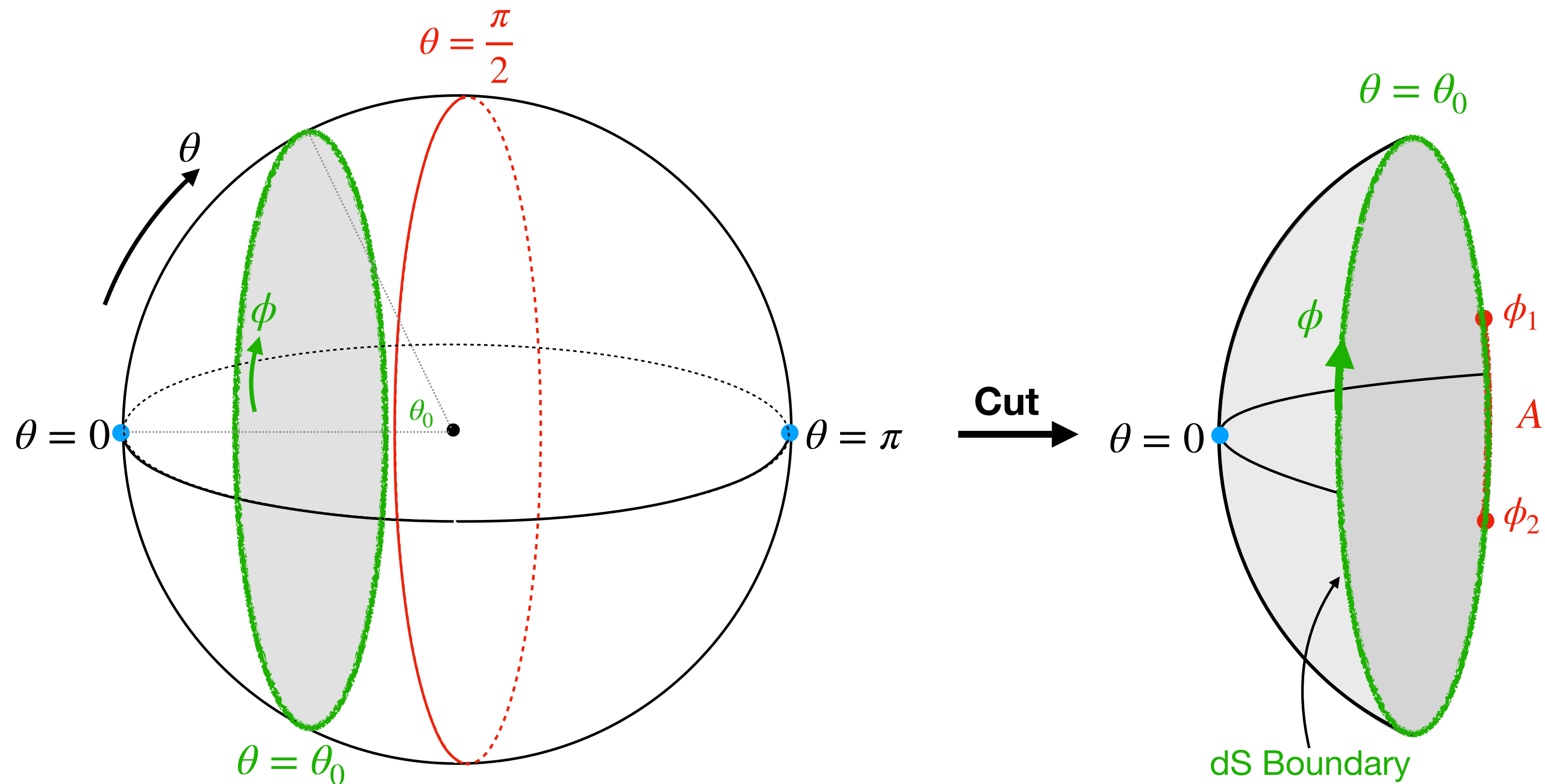
$$S_A^{\text{dis,cyl}} = \frac{c}{6} \log \frac{|f(w_1) - \bar{f}(\bar{w}_1)| |f(w_2) - \bar{f}(\bar{w}_2)|}{\epsilon^2 |f'(w_1)| |f'(w_2)|} + 2S_{\text{bdy}}$$

03. Holography for a Half de Sitter

Topic Today: A half de Sitter space



Topic Today: A half de Sitter space



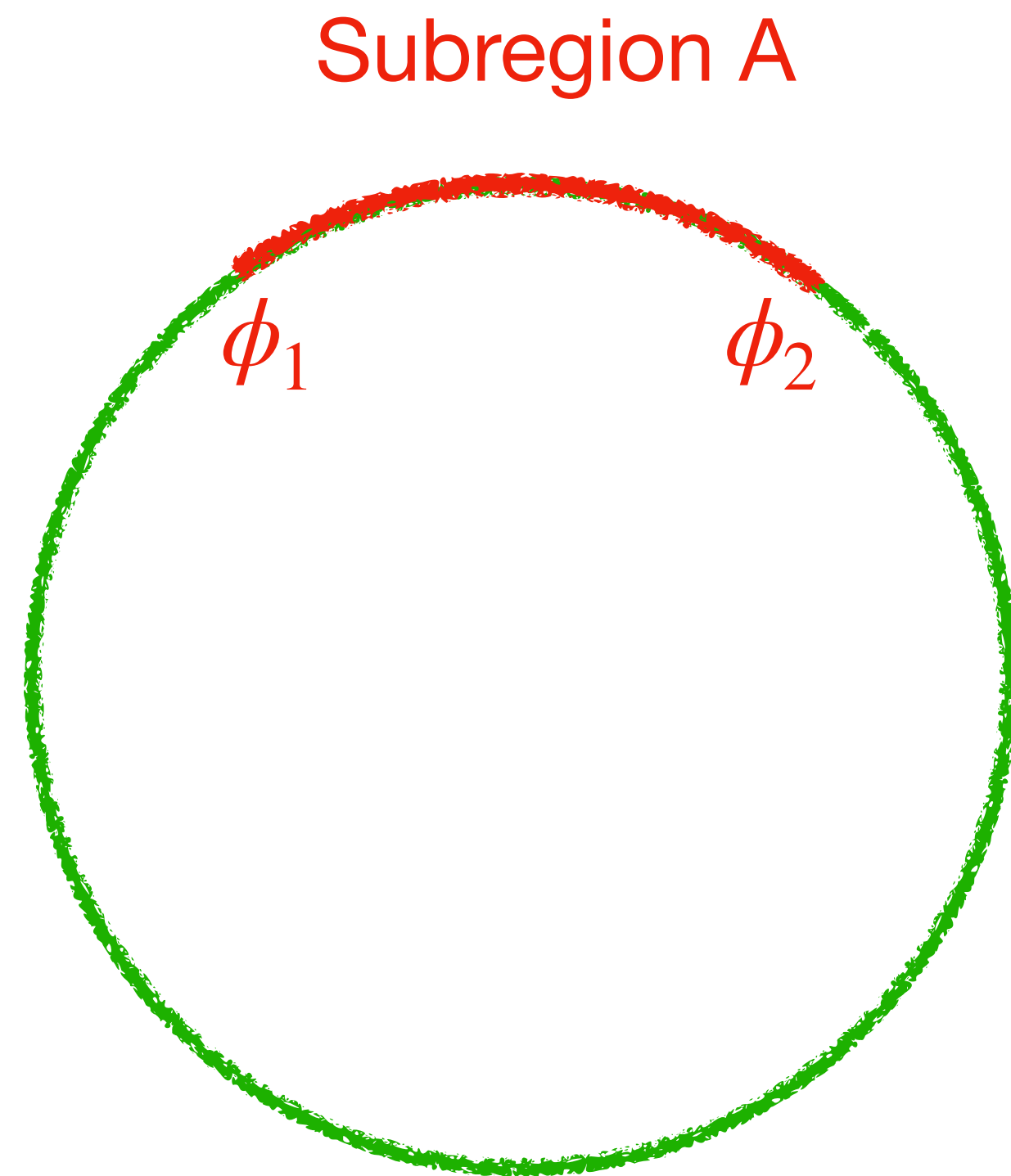
$$ds^2 = -dt^2 + \cosh^2 t (d\theta^2 + \sin^2 \theta d\phi^2)$$

3-dim de Sitter bulk spacetime

$$ds^2 = -dt^2 + \cosh^2 t \sin^2 \theta_0 d\phi^2$$

2-dim de Sitter spacetime

Perspective 01: Holographic Entanglement Entropy Without EOW brane



RT formula in dS

$$S_A = \frac{D_{12}}{4G_N} \quad ?$$

Perspective 01: Holographic Entanglement Entropy Without EOW brane

Bulk spacetime dS_3 $ds^2 = -dt^2 + \cosh^2 t(d\theta^2 + \sin^2 \theta d\phi^2).$

Extremal surface \rightarrow spacelike geodesics

Spacelike geodesics exist only for $\cos^2 \theta_0 + \sin^2 \theta_0 \cos(\phi_1 - \phi_2) \geq 1 - \frac{2}{\cosh^2 t_0}.$

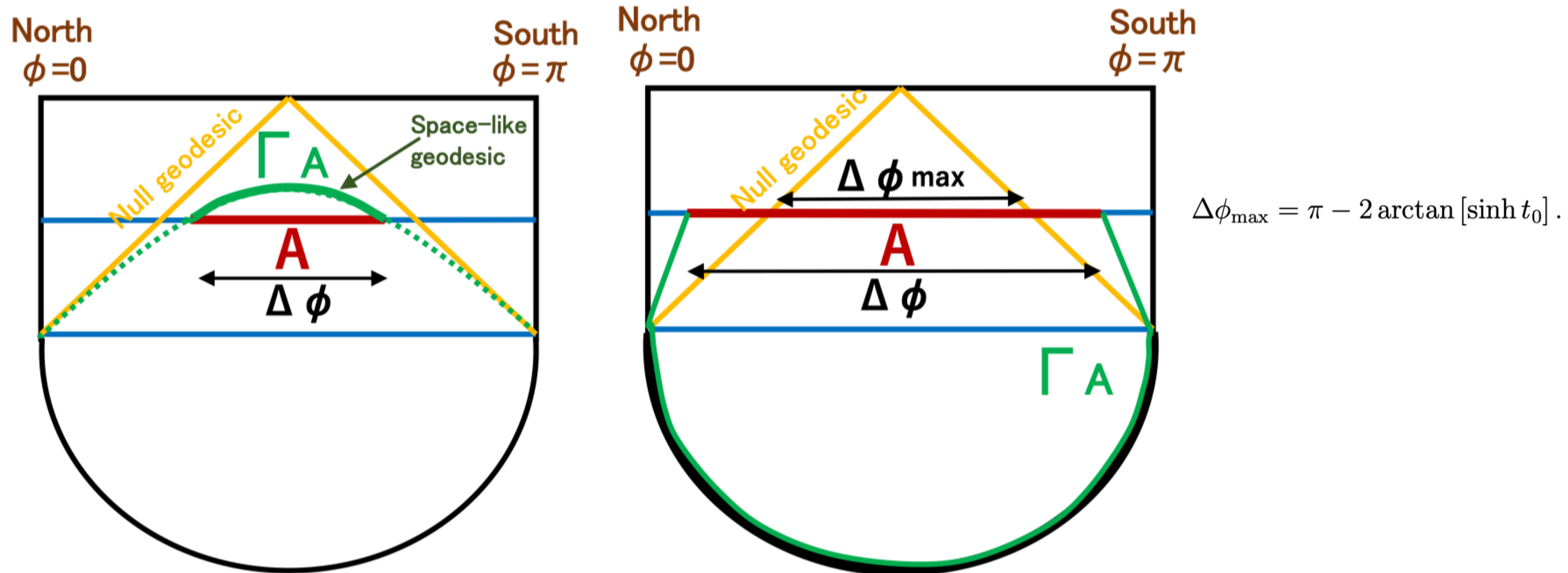
It is always violated at late times for any fixed subregions

Complex Geodesics Length/Area of Minimal Surface

$$D_{12} = \pi + i \operatorname{arccosh} [\sinh^2 t - (\cos^2 \theta_0 + \sin^2 \theta_0 \cos(\phi_1 - \phi_2)) \cosh^2 t_0].$$

Perspective 01: Holographic Entanglement Entropy Without EOW brane

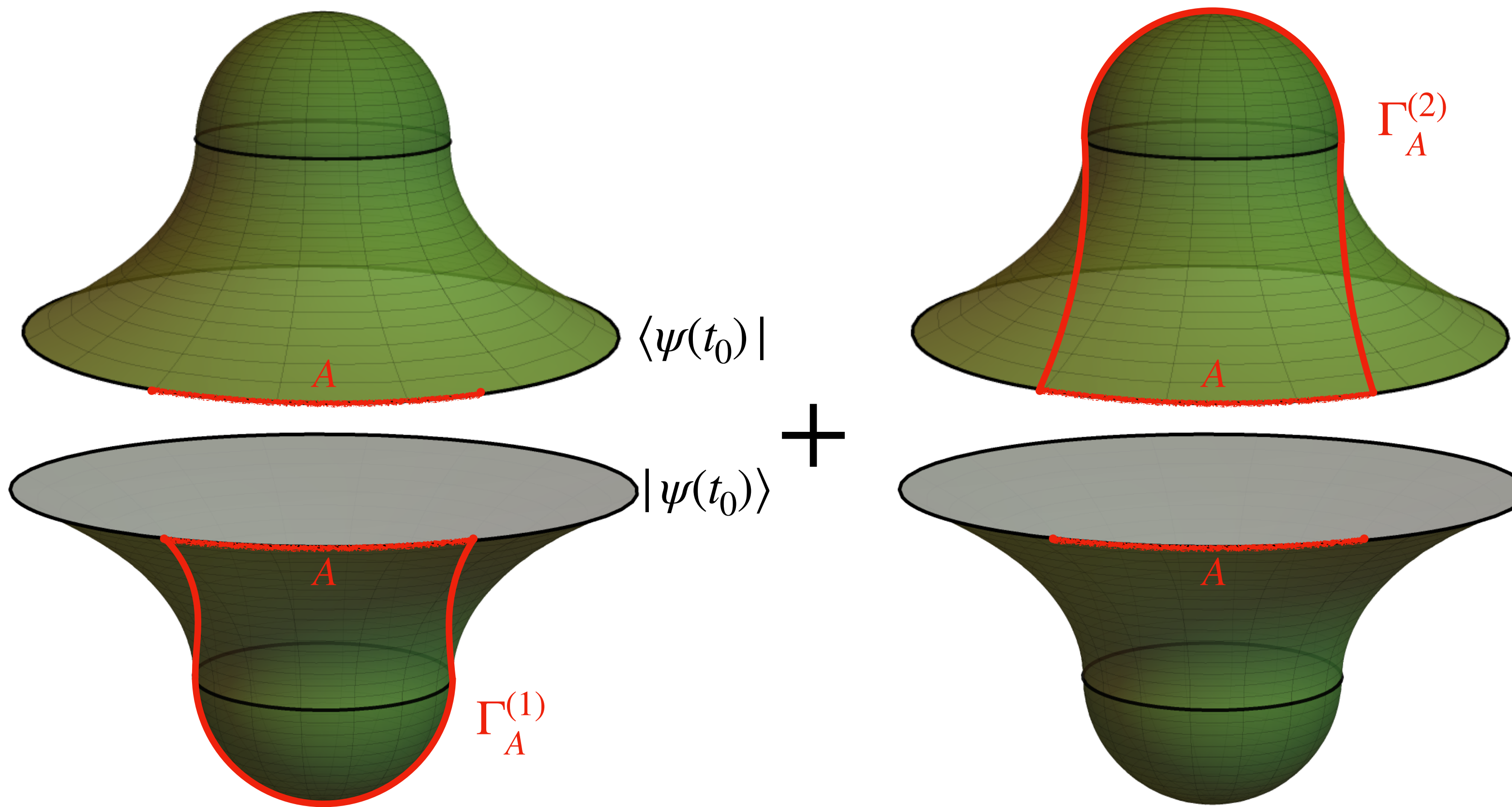
2-dim de Sitter Boundary at $\theta_0 = \pi/2$



$$D_{12} = \pi + i \operatorname{arccosh} \left[\sinh^2 t - (\cos^2 \theta_0 + \sin^2 \theta_0 \cos(\phi_1 - \phi_2)) \cosh^2 t_0 \right].$$

Perspective 01: Holographic Entanglement Entropy Without EOW brane

There are two geodesics: in ket and bra geometry



$$Z_{\text{tot}}^{(n)} = e^{(1-n)S} + e^{(1-n)S^*},$$

$$S_A = \lim_{n \rightarrow 1} \frac{1}{1-n} \log Z_{\text{tot}}^{(n)} \simeq \frac{S + S^*}{2} = \text{Re}[S].$$

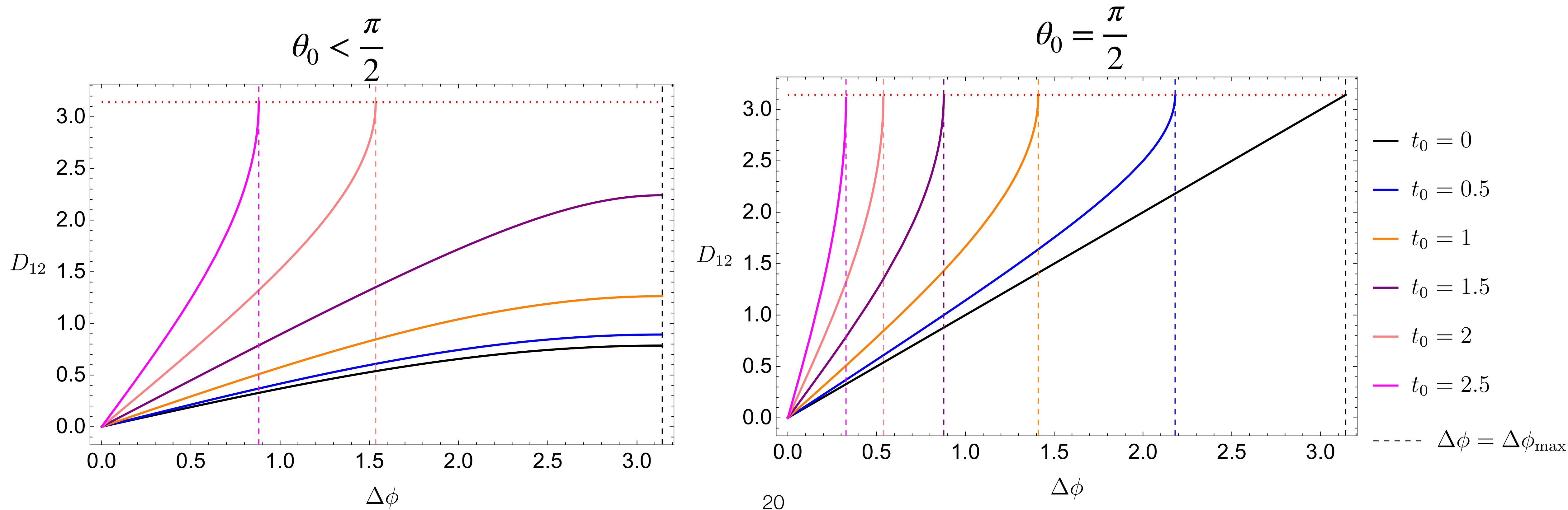
A half dS entropy

$$S_A = \frac{\pi}{4G_N} = \frac{1}{2} S_{\text{dS}},$$

Perspective 01: Holographic Entanglement Entropy Without EOW brane

Spacelike geodesics exist only for $\cos^2 \theta_0 + \sin^2 \theta_0 \cos(\phi_1 - \phi_2) \geq 1 - \frac{2}{\cosh^2 t_0}$.

Geodesic Length at different time slices



Perspective 01: Holographic Entanglement Entropy Without EOW brane

But it is a convex function at late times!

$$\frac{d^2 D_{12}}{d\Delta\phi d\Delta\phi} = \frac{\sqrt{2} \cosh t \sin \theta_0 \sin\left(\frac{\theta_0}{2}\right) (\cosh^2 t \sin^2 \theta_0 - 1)}{(1 + \cos D_{12})^{3/2}} > 0$$

$$t > t_{\max} = \operatorname{arccosh}\left(\frac{1}{\sin \theta_0}\right) = \log\left(\cot\left(\frac{\theta_0}{2}\right)\right). \quad \text{Intersection with cosmological horizon}$$

Violation of (strong) subadditivity!

Perspective 01: Holographic Entanglement Entropy

Subadditivity

$$S_A + S_B \geq S_{AB}$$

Strong Subadditivity (SSA)

$$S_{AB} + S_{BC} \geq S_{ABC} + S_B$$

Violation of (strong) subadditivity!



Boundary theory is highly non-local (but different from $\overline{\text{TT}}$ deformation)

Example: $S = \int dx^d \phi(x) e^{(-\partial_x^2)^q} \phi(x),$ $S_A \propto \left(\frac{L_A}{\epsilon}\right)^{d-2+2q},$ [Li, Takayanagi 2010]

entanglement entropy defined from the replica method (Euclidean path-integral) is not guaranteed to satisfy the subadditivity!

Perspective 01: Holographic Entanglement Entropy

Subadditivity

$$S_A + S_B \geq S_{AB}$$

Strong Subadditivity (SSA)

$$S_{AB} + S_{BC} \geq S_{ABC} + S_B$$

Violation of (strong) subadditivity!

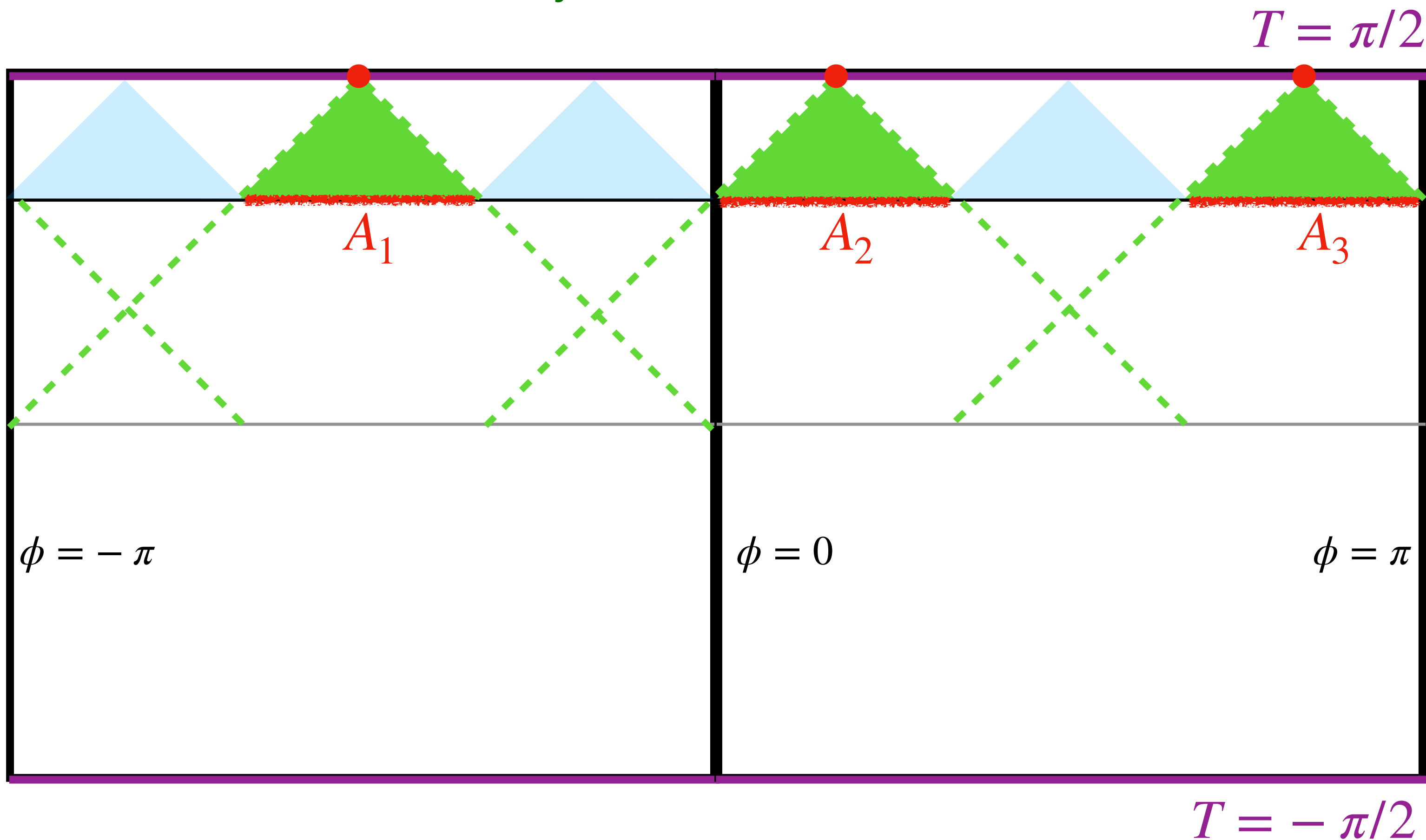


The different time slices are not just a unitary/isometric evolution in dS space

Perspective 01: Holographic Entanglement Entropy

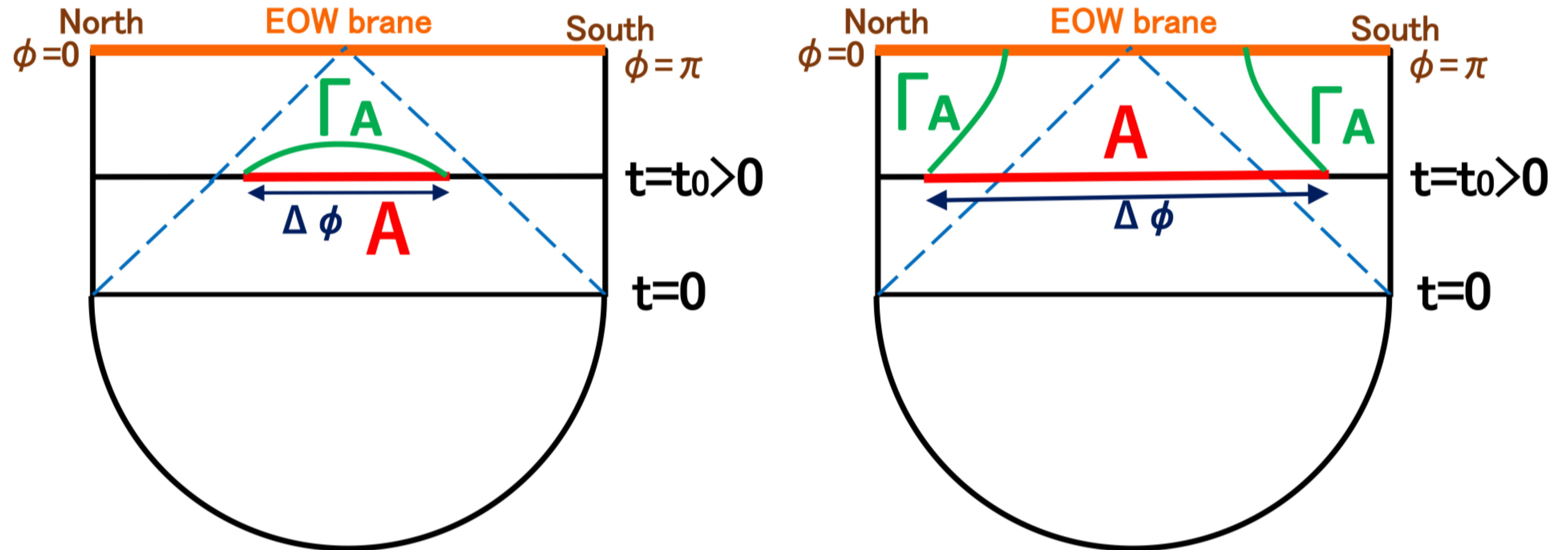
Overcounting Hilbert Space

2-dim de Sitter boundary



$$S_{A_1 \cup A_2 \dots \cup A_N} = N \times \frac{1}{2} S_{\text{dS}} \rightarrow \infty$$

Perspective 02: Holographic Pseudo Entropy with EOW brane



See our paper for more details :
Holographic pseudo entropy, higher dimensional de Sitter, different EOW branes

Thanks for your attention!